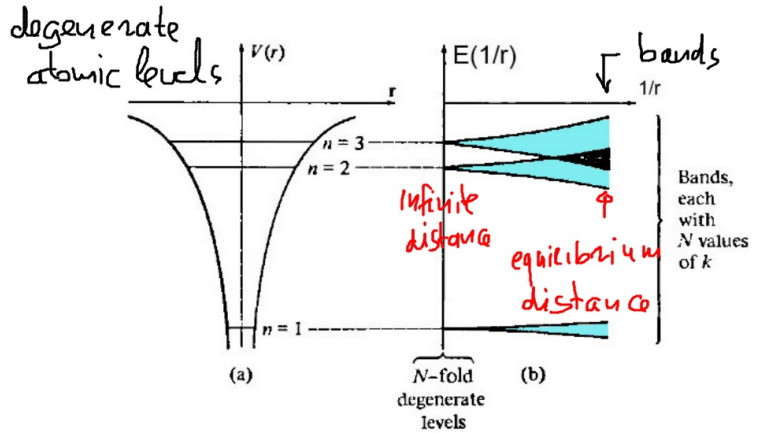


Solid State Physics

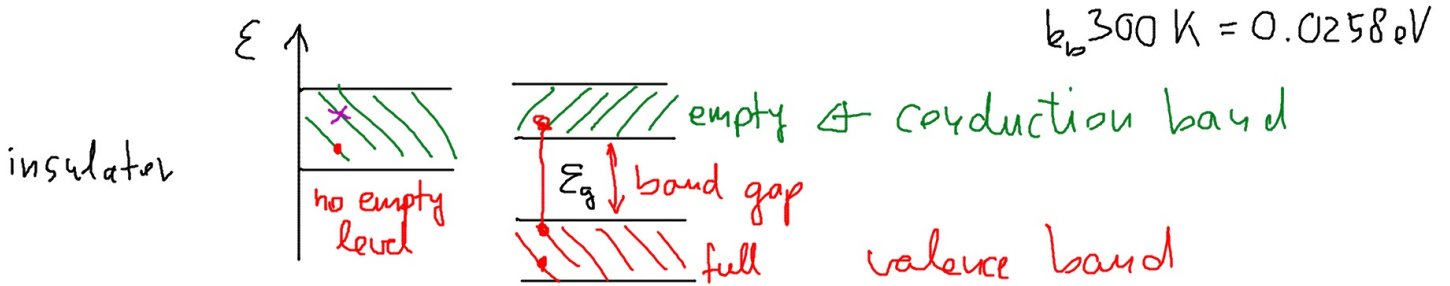
Electrical properties

insulators	metals
no/small/ negligible overlap of electron orbitals	large overlap

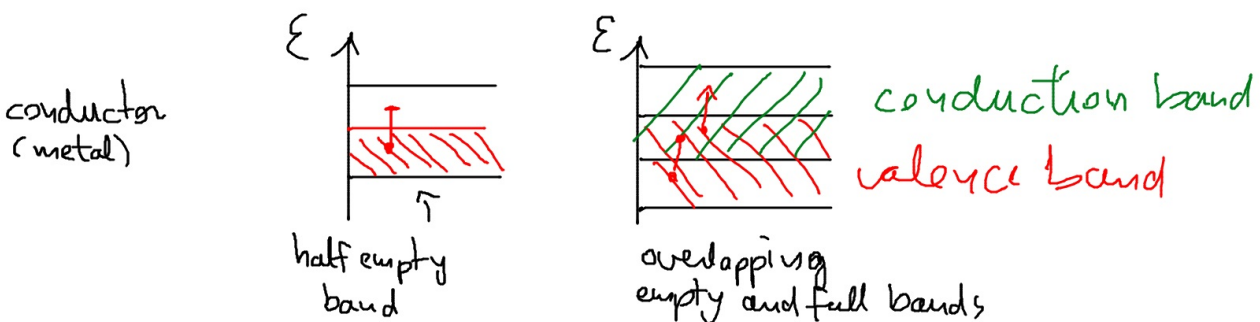
overlap breaks
degeneration
bands form



band from one level
 $\Rightarrow \sim 10^{-20}$ eV



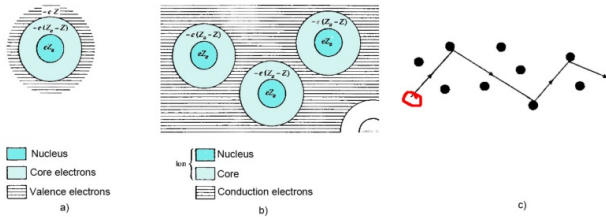
if $E_g \gg k_B T \Rightarrow$ insulator
 if $E_g \sim k_B T \Rightarrow$ semiconductor (also insulator)



classical model of conduction

Drude-model

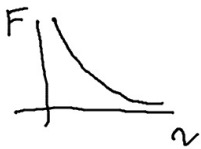
~ 1900



1. ion cores not moving
2. non-interacting electrons (electron-gas)
3. only electron-ion collisions
4. instantaneous
5. $P_{\text{collision}} \sim \Delta t$
6. $v = 0$ after collision

$Q_{\text{electron}} = -e$

$\underline{E} \neq 0 \quad \underline{F} = -e\underline{E} \quad \begin{matrix} \underline{E} \\ \downarrow \\ \underline{a} \end{matrix}$



$a = \frac{F}{m_e} = \frac{eE}{m_e}$

$v(t) = \cancel{v_0} + at \Leftrightarrow dv = a dt$

collision occurs after $\tau \quad v(\tau) = \frac{eE}{m_e} \tau$

not the electron velocity, just the change caused by \underline{E}

after $v(\tau+0) = 0$

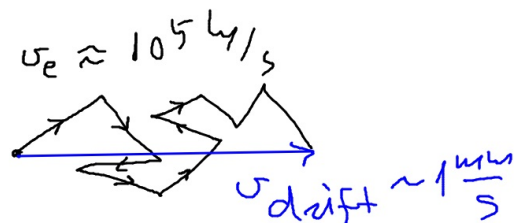
no $\frac{1}{2}$ as τ is average too

$\langle v \rangle = \frac{1}{\tau} v_{\text{max}} = \frac{eE}{m_e} \tau = \mu E$
 \uparrow mobility $\mu = \frac{e\tau}{m_e}$

v_{drift} average collision time

metal
 $n \approx 10^{28} \text{ m}^{-3}$

velocity
 $v_{\text{drift}} = -\frac{eE}{m_e} \tau$



$\underline{j} = -en v_{\text{drift}}$

conductivity

$\underline{j} = -en v_{\text{drift}} = -n \frac{e^2 E}{m_e} \tau = -n \frac{e^2 \tau}{m_e} E = -\sigma \underline{E}$

can be measured $\rightarrow \sigma = \frac{ne^2 \tau}{m_e}$ calculated

$\underline{j} = -en v_{\text{drift}} = -ne\mu E$

$\sigma = ne\mu$

Silver $\tau \approx 4 \times 10^{-14} \text{ s}$

units $[\sigma] = \frac{[j]}{[E]} = \frac{A/m^2}{V/m} = \frac{1}{\Omega m}$ resistivity $[\rho] = \Omega m$
 or $\Omega \frac{mm^2}{m}$

$[\mu] = \frac{[G_{drift}]}{[E]} = \frac{m/s}{V/m} = \frac{m^2}{Vs}$

calculate n

$n = L_A \frac{Z \rho_m}{A}$

of conduction electrons/atom
atomic mass in kg/mol
mass density kg/m³

example: Al
 $\rho_m = 2700 \frac{kg}{m^3}$
 $Z = 3$
 $A = 0.02698 \frac{kg}{mol}$

$n = 6.024 \cdot 10^{23} \frac{atoms}{mol} \cdot \frac{3 \cdot 2700 \frac{kg}{m^3}}{0.02698 \frac{kg}{mol}}$
 $= 1.807 \cdot 10^{29} \frac{electron}{m^3}$

τ & calculations hard

using σ : $\sigma = \frac{ne^2}{m} \tau \rightarrow \tau = \frac{m \sigma}{ne^2} = \frac{m}{\rho n e^2}$

$\tau = \frac{m}{L_A \frac{Z \rho_m}{A} e^2 \cdot \rho} = \frac{m \cdot A}{Z \cdot \rho_m e^2 \cdot L_A}$

example: silver (Ag)

$\rho(0K) = 1.51 \cdot 10^{-8} \Omega m$ ($\rho = \rho(T)$)

$Z = 1$ (5s¹)

$A = 107.8682 \text{ g/mol}$

$\rho_m = 10.49 \text{ g/cm}^3$

$\tau = 4.013 \cdot 10^{-14} \text{ s}$

mean free path \bar{l}

$\tau = \frac{\bar{l}}{\langle v \rangle}$

Drude: $\langle v \rangle = v_{th} = \sqrt{\frac{3k_B T}{m_e}}$

$\sigma = \frac{ne^2 \tau}{m_e} \rightarrow \rho = \frac{m_e}{ne^2 \tau} \sim \sqrt{T}$

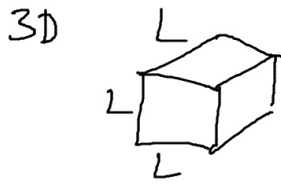
reality: $\rho \sim T$

Failures of the Drude model:

$\rho(T), C_V, \frac{N}{\sigma}$

Spurmerfeld model

1. cond. e does not interact w. by cores or e.
2. potential box
3. Pauli principle true



N electrons

$$V(x, y, z) = \begin{cases} 0 & 0 \leq x \leq L \text{ and } 0 \leq y \leq L, 0 \leq z \leq L \\ \infty & x > L \text{ or } y > L \text{ or } z > L \\ & x < 0 \quad y < 0 \quad z < 0 \end{cases}$$

$$\Psi(\underline{n}_1, \underline{n}_2, \dots, \underline{n}_N) = \sum_{i=1}^N \Phi_i(\underline{n}_i)$$

$$\Phi_i(\underline{n}_i) \equiv \underbrace{\Phi_{i_x}(x)}_{\Phi(x)} \cdot \underbrace{\Phi_{i_y}(y)}_{\Phi(y)} \cdot \underbrace{\Phi_{i_z}(z)}_{\Phi(z)}$$

$$\Phi(x) = \sqrt{\frac{2}{L}} \sin k_x x \quad |k_x = \frac{\pi}{L} \cdot n_x, n_x = 1, 2, 3$$

$$\Phi(y) = \sqrt{\frac{2}{L}} \sin k_y y \quad |k_y = \frac{\pi}{L} \cdot n_y, n_y = 1, 2, 3$$

$$\Phi(z) = \sqrt{\frac{2}{L}} \sin k_z z \quad |k_z = \frac{\pi}{L} \cdot n_z, n_z = 1, 2, 3$$

1D a-lattice const $L = N \cdot a$, N - atoms $\Delta k = \frac{\pi}{L} \sim 10^{-13} - 10^{-15} \frac{1}{m}$

$$\Sigma_v(L) = \frac{\hbar^2 k^2}{2m_e} = \frac{\hbar^2 \left(\frac{\pi}{L}\right)^2 n^2}{2m_e} = \frac{\hbar^2 \pi^2}{2m_e N a^2} \cdot n^2$$

3D $\Psi(x, y, z) = \left(\frac{2}{L}\right)^{3/2} \sin k_x x \cdot \sin k_y y \cdot \sin k_z z$

$$\Sigma = \frac{\hbar^2 k^2}{2m_e} = \frac{\hbar^2}{2m_e} \cdot (k_x^2 + k_y^2 + k_z^2)$$

ϵ_F

$$\Delta V_k = \left(\frac{2\pi}{L}\right) \left(\frac{2\pi}{L}\right) \left(\frac{2\pi}{L}\right) = \left(\frac{2\pi}{L}\right)^3 = \frac{8\pi^3}{V}$$

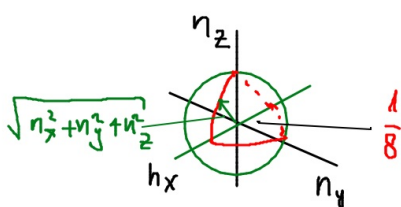
of k points in Ω_k : $n = \frac{\Omega_k}{\Delta V_k} = \frac{V}{\pi^3} \cdot \Omega_k$

density of levels =

$$\rho(\text{levels}) = \frac{n}{\Omega_k} = \frac{V}{\pi^3}$$

for N k-states ($N \gg 1$) $\Rightarrow \Omega(N)$ is a sphere

$$\Omega(N) = \frac{4}{3} k_F^3 \pi \quad \text{but } n_x, n_y, n_z > 0 \Rightarrow \frac{1}{8} \times \Omega(N)$$



$$N_{\text{states}} = \text{# of levels} \cdot \frac{\Omega(\mathcal{N})}{8}$$

$$N_{\text{states}} = \frac{V}{\pi^2} \cdot \frac{1}{8} \cdot \frac{4}{3} k_F^3 \pi = \frac{V}{6\pi^2} k_F^3$$

2 electrons/atom $N_{\text{electrons}} = Z \cdot N$

Pauli $\rightarrow N_{\text{states}} = \frac{N_{\text{electrons}}}{2}$

$$\frac{V}{6\pi^2} k_F^3 = \frac{N_{\text{electrons}}}{2}$$

$$N_e = N_{\text{electrons}}$$

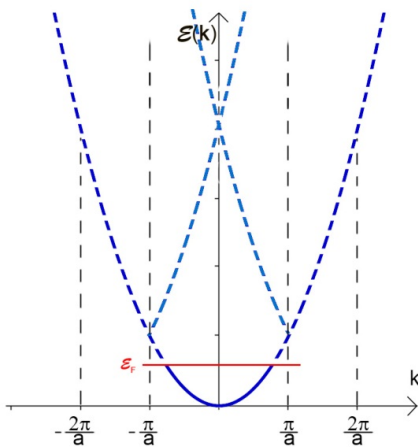
$$N_e = \frac{V}{3\pi^2} k_F^3 \quad n = \frac{N_e}{V} = \frac{k_F^3}{3\pi^2}$$

$$k_F = \frac{2\pi}{a} \Rightarrow k \leq k_F \text{ value, } k \in [0, \frac{2\pi}{a}]$$

$$\epsilon_F = \frac{\hbar^2 k_F^2}{2m} \rightarrow k_F = \sqrt{\frac{2m\epsilon_F}{\hbar^2}}$$

$$v_F = \frac{\hbar k_F}{m_e}$$

$$n = \frac{1}{3\pi^2} \left(\frac{2m\epsilon_F}{\hbar^2} \right)^{3/2} = \frac{(2m\epsilon_F)^{3/2}}{3\pi^2 \hbar^3}$$



$$(k_F) \sim 0.1 \text{ nm}^{-1} \checkmark$$

$$v_F \sim 0.001 c \quad \times$$

$\epsilon_F \sim$ binding energy

example

Cu $\epsilon_F = 4.1 \text{ eV}$

$$n = \frac{(2m\epsilon_F)^{3/2}}{3\pi^2 \hbar^3} = \frac{(2 \cdot 9.1 \cdot 10^{-31} \text{ kg} \cdot 4.1 \text{ eV} \cdot 1.6 \cdot 10^{-19} \text{ J/eV})^{3/2}}{3\pi^2 (1.57 \cdot 10^{-34})^3}$$

$$= 3.77 \cdot 10^{28} \frac{\text{electron}}{\text{m}^3} \quad \leftarrow \text{wrong}$$

$$v_F = 1.2 \cdot 10^6 \text{ m/s} = 0.004 c$$

Specific heat of metals # of electrons in $\Delta\epsilon$

$$\Delta n_{\text{states}} = g(\epsilon) \Delta\epsilon$$

$$g(\epsilon) = \frac{8\pi (2m\epsilon)^{1/2}}{\hbar^3} \cdot \sqrt{\epsilon}$$

$$f_{F-D} = \frac{1}{e^{(\epsilon - \epsilon_F)/k_B T} + 1}$$

$$dn = f_{F-D}(\epsilon) g(\epsilon) d\epsilon$$

$$U = \int_0^{\infty} \epsilon dn = \int_0^{\infty} \epsilon \cdot g(\epsilon) f_{F-D}(\epsilon) d\epsilon$$

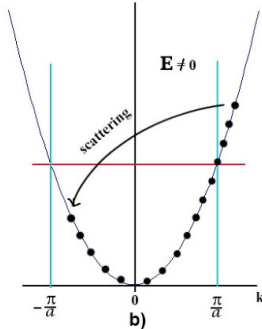
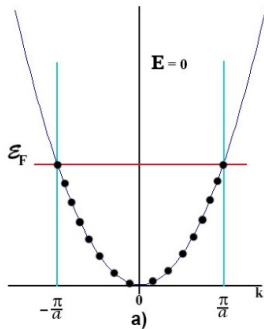
$$U = \int_0^{\infty} \frac{\epsilon \cdot g(\epsilon)}{e^{(\epsilon - \epsilon_F)/k_B T} + 1} d\epsilon =$$

$$= \int_0^{\epsilon_F - k_B T} (\) d\epsilon + \int_{\epsilon_F - k_B T}^{\epsilon_F + k_B T} (\) d\epsilon + \int_{\epsilon + k_B T}^{\infty} (\) d\epsilon$$

\downarrow $\sim \text{const}$
 \downarrow ~ 0

$$C_V = \frac{\pi^2}{3} k_B^2 g(\epsilon_F) T$$

$$U = \text{const} + \int \frac{g(\epsilon) \epsilon}{(\)} d\epsilon$$



$$\approx g(\epsilon_F) \cdot \int_{\epsilon_F - k_B T}^{\epsilon_F + k_B T} \frac{\epsilon}{e^{(\epsilon - \epsilon_F)/k_B T} + 1} d\epsilon$$

$$v \sim v_F$$

$$\bar{v} = \frac{\bar{v}}{v_F}, \quad \bar{v} = \frac{n e^2 \tau}{m_e} = \frac{v_F^2}{M R} \frac{\bar{v}}{v_F}$$

$$v_F \neq v_F(T) \quad \bar{v} \sim \frac{1}{T} \Rightarrow \rho \sim T \quad \checkmark$$

Problems

- $v_F \sim 10 \times v_F^{(Duda)}$ \times
- $\bar{v} = 10^{-7} \text{ m/s}$ \times $\bar{v} \sim 10^{-9} \text{ m/s}$
- insulators ?
- Hall effect
- magnetoresistance